# **DS10:** Potentials and the Method of Images

Solution set

Philip Cherian 29th October, 2018

Ashoka University

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The Multipole Expansion

We have seen that:

$$V_{\text{dipole}} = \frac{p\cos\theta}{r^2} = \frac{\mathbf{p}.\hat{\mathbf{r}}}{r^2} \tag{1}$$

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We also know that  $\mathbf{E} = -\nabla V$ , and so:

$$\mathbf{E}_{\text{dipole}} = -\left(\frac{\partial V}{\partial r}\hat{r} + \frac{1}{r}\frac{\partial V}{\partial \theta}\hat{\theta}\right)$$

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Now use the fact that  $\mathbf{p} = (\mathbf{p} \cdot \hat{r})\hat{r} + (\mathbf{p} \cdot \hat{\theta})\hat{\theta} = p\cos\theta\hat{r} - p\sin\theta\hat{\theta}$  and the answer should follow. (See Griffiths' Electric Field of a Dipole – pg 153 in my edition, and Problem 3.33).

2

**Charge Distributions** 

We will be using the differential form of Gauss's law:  $\nabla \cdot \mathbf{E} = \rho/\epsilon_0$ . Thus, the charge distribution is given by

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Computing the divergence in spherical polar coordinates:

$$\rho(r,\theta,\phi) = \epsilon_0 \left( \frac{1}{r^2} \frac{\partial}{\partial r} \left( r^2 \frac{A}{r} \right) + \frac{1}{r \sin \theta} \frac{\partial}{\partial \theta} \left( 0 \right) + \frac{1}{r \sin \theta} \frac{\partial}{\partial \phi} \left( \frac{B \sin \theta \cos \phi}{r} \right) \right)$$

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$$= \epsilon_0 \left( \frac{A - B \sin \phi}{r^2} \right)$$
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The electric field is usually assumed to go to zero far away from the charges. Or, if this isn't given to us, we sometimes use symmetry to guess the form of the field (for example, in the case of infinite sheets, etc.).

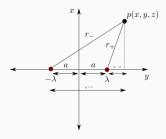
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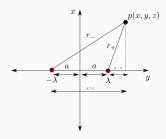
The electric field is usually assumed to go to zero far away from the charges. Or, if this isn't given to us, we sometimes use symmetry to guess the form of the field (for example, in the case of infinite sheets, etc.). However, in this case, the distribution is infinite with no symmetries we can exploit. In fact, it is a badly posed problem: there can be no answer.

(Try calculating the field using Gauss's law, and you'll get an integral that diverges.)

**Equipotential Surfaces** 



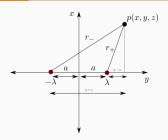
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Thus, we have that:

$$V(x, y, z) = \frac{\lambda}{4\pi\epsilon_0} \ln\left(\frac{(y+a)^2 + x^2}{(y-a)^2 + x^2}\right)$$
 (6)

-

Let us set the potential to be constant  $(V_0)$ . Then, taking the exponential on both sides in the previous equation,

$$\left(\frac{(y+a)^2 + x^2}{(y-a)^2 + x^2}\right) = e^{\frac{4\pi\epsilon_0 V_0}{\lambda}} = C$$
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$$x^{2} + y^{2} + a^{2} - 2ay\left(\frac{C+1}{C-1}\right) = 0$$

$$x^{2} + (y-y_{0})^{2} = R^{2} \implies x^{2} + y^{2} + y_{0}^{2} - R^{2} - 2yy_{0} = 0$$
(8)

Comparing, we see that these are circles centered at  $y_0$ , with radius R.

$$y_0 = a\left(\frac{C+1}{C-1}\right)$$

$$R = 2a\frac{\sqrt{C}}{|C-1|}$$
(9)

Exercise: Show that in terms of  $V_0$  this is

$$y_0 = a \coth\left(\frac{2\pi\epsilon_0 V_0}{\lambda}\right)$$

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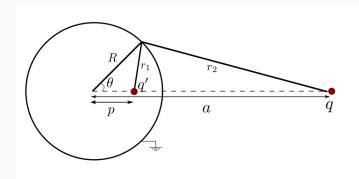
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- 2. When  $V_0 \to 0$  Both the centre and radius of the circle go to  $\infty$ ! But wait, that's just the x-axis, our reference!  $\implies$  No problem.

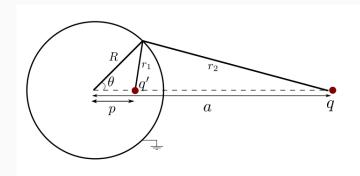
The Method of Images

#### Questions 1,2, and 3



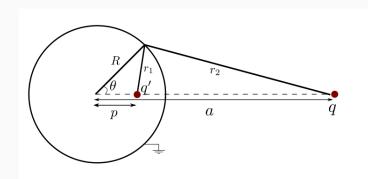
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- 1. The charge on the conductor will rearrange itself so that any changes occurring inside the shell will not affect the potential outside it.
- 2. Grounded  $\implies$  the potential on the shell is the same as at infinity: zero.
- 3. If the potential is zero on the surface, then

$$\frac{1}{4\pi\epsilon_0} \left( \frac{q}{r_2} + \frac{q'}{r_1} \right) = 0 \implies \frac{q'}{r_1} = -\frac{q}{r_2} \tag{11}$$

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- 2. In terms of the law of cosines, we can write

$$r_1^2 = R^2 + p^2 - 2Rp\cos\theta$$
  

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(12)

from which the equation follows.

3. What is the induced charge on the shell?

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- 4. The induced charge on the must be the same as the image charge.
- 5. But the induced charge on the shell is negative if *q* is positive. Thus, so is the image charge.

$$q'^2R^2 + q'^2a^2 - 2q'^2Ra\cos\theta = q^2R^2 + q^2p^2 - 2q^2Rp\cos\theta$$

$$q^{\prime 2}R^2 + q^{\prime 2}a^2 - 2q^{\prime 2}Ra\cos\theta = q^2R^2 + q^2p^2 - 2q^2Rp\cos\theta$$
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Thus, we get two equations:

$$q^{2} = q^{2} \frac{a}{p} \implies q' = -q \sqrt{\frac{p}{a}}$$

$$\frac{p}{a} (R^{2} + a^{2}) = (R^{2} + p^{2}) \implies p = \frac{R^{2} + a^{2}}{2a} \pm \frac{1}{2} \sqrt{\frac{(R^{2} + a^{2})^{2}}{a^{2}} - 4\frac{R^{2}a^{2}}{a^{2}}}$$
(14)

#### **Solution**:

$$q^{\prime 2}R^2 + q^{\prime 2}a^2 - 2q^{\prime 2}Ra\cos\theta = q^2R^2 + q^2p^2 - 2q^2Rp\cos\theta$$
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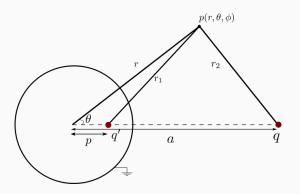
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Solution:

$$q' = -q\frac{R}{a}$$

$$p = \frac{R^2}{a} \quad \text{or} \quad p = a$$
(15)



1. Consider an arbitrary point outside the shell  $p(r,\theta,\phi)$ . To find the potential at p, we use the same technique as before, except we replace  $R \to r$ .

## Question 6, 7

$$V(p) = \frac{1}{4\pi\epsilon_0} \left( \frac{q}{r_2} + \frac{q'}{r_1} \right) = \left( \frac{q}{\sqrt{r^2 + a^2 - 2ra\cos\theta}} + \frac{-q\frac{R}{a}}{\sqrt{r^2 + p^2 - 2rp\cos\theta}} \right)$$
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- 2. Thus, same problem, except with two image charges:

$$q' = -q\frac{R}{a}$$

$$q'' = 4\pi\epsilon_0 RV_0$$
(17)

$$V_{\text{ungrounded}}(p) = V_{\text{grounded}}(p) + \frac{q''}{4\pi\epsilon_0 r}$$

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